

TABLE I. SOLID MATERIAL PROPERTIES

Properties	PDMS	PMMA
Density (kg/m ³)	0.97	1180
Young's Modulus, E (Pa)	3.6x10 ⁵	2.45x
Poisson's Ratio	0.49	0.35

TABLE II. FLUID PROPERTIES

Fluids	Density (kg/m ³)	Dynamic Viscosity (Kg/ms)	
Blood	1050	3.5x1	
Water	997	8.9x1	
Acetone	784	3.16x	
Propanol	803	1.92x	
Acetic Acid	1050	1.15x	
Non Newtonian Fluids		High Shear Viscosity	Low Shear
Blood 35%H**	1050	0.00313	0.
Blood 40%H**	1050	0.0035	0.
Blood 48%H**	1050	0.005	0.

** As Non-Newtonian (Bird-Carreau Model) [17]

III. NUMERICAL ANALYSIS

In order to examine the sensitivity of the two biosensors, a computational study was performed. Both the fluid flow and the fluid-structure interaction were included in the computational model. A mesh-independency analysis was performed for this study and the optimum element size was found.

A. Boundary Conditions.

For both fluid domains and biosensors the boundary conditions are the same. The fluid has constant inlet velocity 0.1 mm/s and zero pressure is applied at the outlet. The no slip/no penetration conditions are also applied on the walls. The boundary conditions for the Biosensor B1 are presented in (Fig. 2). The boundary conditions are exactly the same with the B2 design. The fluid flow was considered laminar. Gravity was taken into account for the Computational Fluid Dynamics (CFD) and Mechanical Analysis.

B. Mathematical Model

Using ANSYS we performed the analysis of the fluid flow and the deflection of the microcantilever's beam. The current problem dictates that a fluid flows inside the microfluidics channel and forces the beam to deflect downwards.

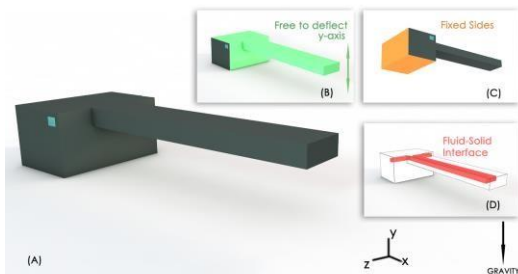


Figure 2. The boundary conditions for the biosensor B1. With green color are presented the faces which are free to deform. The orange color depicts the fixed sides of the structure, while the red color shows the fluid-structure interface.

The mechanical analysis for the deflection of the suspended microcantilever's beam was employed using the following system of equations

$$[K]\{u\} = \{F\}, \quad (1)$$

where, K is the stiffness matrix of the material, u, is the nodal displacement vectors and F, is the external loads. For the CFD analysis, three components of the velocity field $V = v \hat{i} + v \hat{j} + v \hat{k}$ and the pressure field within the $x \ y \ z$ control volume, have to be determined. Thus, using conservation principles, the mass and the momentum conservation equations are given by the Eqs. (2) and (3a, 3b and 3c), respectively. where, ρ is the fluid density, P is the pressure, μ is the dynamic

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho v_x) + \frac{\partial}{\partial y}(\rho v_y) + \frac{\partial}{\partial z}(\rho v_z) = 0, \quad (2)$$

$$\rho \left(\frac{\partial v_x}{\partial t} + v_x \frac{\partial v_x}{\partial x} + v_y \frac{\partial v_x}{\partial y} + v_z \frac{\partial v_x}{\partial z} \right) = -\frac{\partial P}{\partial x} - \mu \left(\frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial z} \right) + \rho g_x, \quad (3a)$$

$$\rho \left(\frac{\partial v_y}{\partial t} + v_x \frac{\partial v_y}{\partial x} + v_y \frac{\partial v_y}{\partial y} + v_z \frac{\partial v_y}{\partial z} \right) = -\frac{\partial P}{\partial y} - \mu \left(\frac{\partial \tau_{yx}}{\partial x} + \frac{\partial \tau_{yy}}{\partial y} + \frac{\partial \tau_{yz}}{\partial z} \right) + \rho g_y, \quad (3b)$$

$$\rho \left(\frac{\partial v_z}{\partial t} + v_x \frac{\partial v_z}{\partial x} + v_y \frac{\partial v_z}{\partial y} + v_z \frac{\partial v_z}{\partial z} \right) = -\frac{\partial P}{\partial z} - \mu \left(\frac{\partial \tau_{zx}}{\partial x} + \frac{\partial \tau_{zy}}{\partial y} + \frac{\partial \tau_{zz}}{\partial z} \right) + \rho g_z, \quad (3c)$$

fluid viscosity, τ is the stress tensor and g is the earth gravity. The dynamic viscosity of a Newtonian fluid is given by:

$$\mu = \frac{\tau}{\dot{\gamma}}, \quad (4)$$

where, $\dot{\gamma}$ is the shear rate. In case of the Non-Newtonian fluid, the viscosity was calculated using the Bird- Carreau model [17]:

$$\mu = \mu_\infty + \frac{(\mu_0 - \mu_\infty)}{(1 + (\lambda \dot{\gamma})^2)^{\frac{1-n}{2}}}, \quad (5)$$

where, μ_0 is the low shear viscosity equal to 0,056 (Pa *s), μ_∞ is the high shear viscosity and it is equal to the value of the Newtonian model, λ is a time constant equal to 3,313 s, n is the power of law index equal to 0,3568 and $\dot{\gamma}$ is the shear rate. For a steady state problem as this one, the transient term $\partial \rho / \partial t$ can be neglected. For the linear elastic materials, the strain, ϵ and the stress, σ need to be determined as:

$$\sigma = E \epsilon, \quad (6)$$

where, E is the Young's modulus. Both the one and the two way FSI were tested. In the case of the one-way FSI, the CFD or the mechanical analysis was employed first and then, the results from the solved problem are used as initial conditions for the second one. The simulation stops when the second problem converges. For the two-way FSI, both simulations are

performed at the same time by exchanging data and trying to reach an overall convergence.

IV. RESULTS

The mechanical behavior of the biosensors with respect to the beam's deflection was tested. First, a comparison between the one and the two way FSI was performed using water, having velocity 0.1 mm/s at the inlet. The results have shown that the biosensor deflects downwards and the deflection values for the one and the two way FSI were quite close, 2.73 and 2.77 μm , respectively. Although, due to the computational time and memory requirement of the two way FSI, the one way was preferred over the two way. The Biosensor B1 exhibits larger deflections, 5.53 μm , than the B2 which deflects 0.94 μm . Thus, only the B1 version was used for the rest of the study. Two polymeric materials were tested, PDMS and PMMA. The deflection of the PDMS and PMMA was found 5.53 and 0.93 μm , respectively. The PDMS is the most suitable material since it deflects more. Three different microcantilever's beams were tested. The simulations demonstrate that the longer the beam the more it deflects. Thus, the sensitivity of the biosensor increases when the beam's length increases. As a disadvantage, the increased computational time is required as a result of the increased number of elements used in the Finite Elements (FE) calculations. Thus, the 420 μm beam length was an acceptable compromise for the current study. Table III shows the beam's length, the number of the elements used in FE computations and its downwards deflection. The beam's length and the deflection are linearly related as it is shown in (Fig. 3). Furthermore, the sensitivity of the biosensor was found to be closely related to the fluid's velocity. The higher the velocity of the fluid at the inlet, the more the biosensor deflects. A linear relationship between these two parameters was found (Fig. 4). Table IV presents the beam's deflection due to the fluid's velocity. Additionally, a linear relationship between the beam's deflection and the fluids dynamic viscosity, exists for both Newtonian and non-Newtonian fluids. The deflection of the Newtonian and Non-Newtonian fluids, with 0.1 mm/s constant inlet velocity, is presented in (Fig. 5). The velocity streamlines of the fluid flow of the B1 biosensor version, as well as the velocity field distribution in the first longest microchannel and the normalized velocity profile alongside the channel's width and height are presented in (Fig. 6). The von Mises stress was calculated and compared with the tensile strength of the material which is 2.24 MPa [18]. The resulting maximum von Mises stress, at the contact point between the biosensor's base and the beam, is 0.00267 MPa, indicating a safe structure. The very end of the tip is deflected more than the other parts of the beam, while the maximum von Mises stress appears on the opposite direction (Fig.7).

TABLE III. BEAM LENGTH VS DEFLECTION

Ca	Length (μm)	No. Elements	Deflection (μm)
1	280	966,972	0.680
2	420	1,336,553	5.534
3	630	6,499,618	13.850

TABLE IV. FLUID VELOCITY VS DEFLECTION

Case	Velocity (mm/s)	Deflection (μm)
1	0 (not moving)	0.942
2	0.1	5.534
3	0.4	22.078
4	0.8	44.272
5	1	55.340

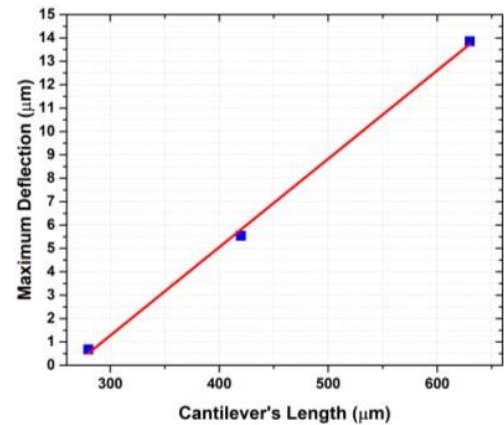


Figure 3. Maximum Deflection of the beam vs Lengths tested.

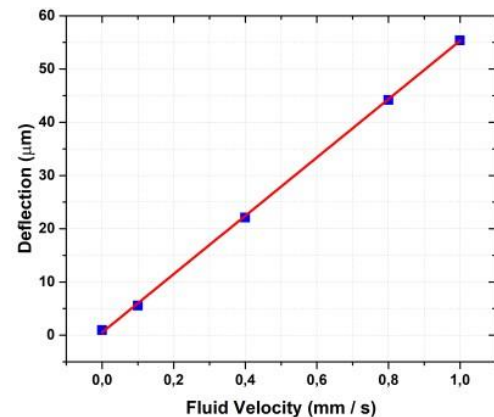


Figure 4. Maximum Deflection of the beam vs Fluid Velocity.

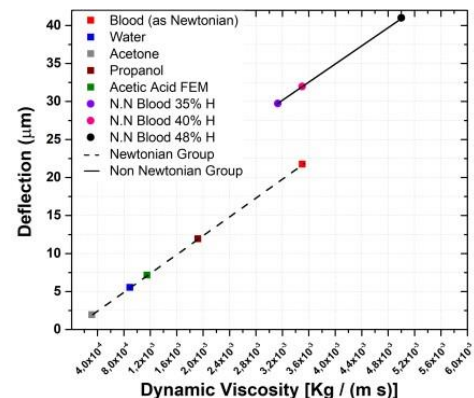


Figure 5. Maximum Deflection of the beam vs Fluid Velocity.

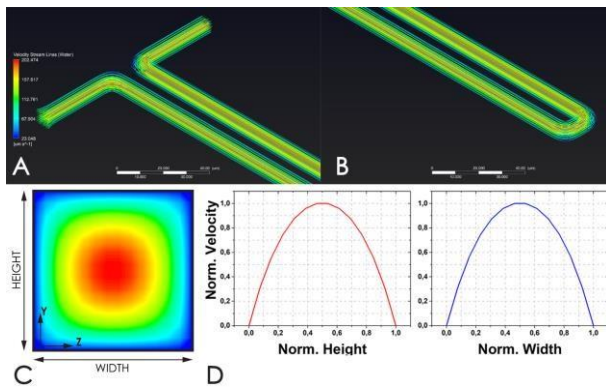


Figure 6. (A) and (B) are the velocity streamlines using water in B1 biosensor version. (C) Velocity field distribution in a cross-section inside the long microfluidics channel. (D) Normalized velocity profiles alongside the channel's height and width.

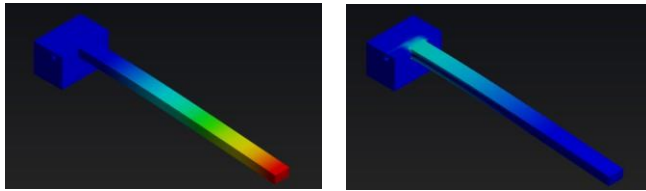


Figure 7. Areas of (Right) Maximum Deflection and (Left) Maximum von Mises Stress of the beam.

IV. DISCUSSION AND CONCLUSIONS

The current biosensor with this small size presents unique and very promising results. The proposed biosensor B1, is able to determine the fluids' dynamic viscosity by measuring the cantilever's beam deflection with high sensitivity. The major advantage is that its encouraging results lead to a more sensitive biosensor than that proposed in the literature [14]. Fig. 8 provides a comparison of the results of this study and the literature [14] in which a bigger geometry, similar to the current B2 biosensor version, with a cantilever size of 6000 x 2000 x 600 (length, width and thickness in μm) and microfluidics cross-section of 200 x 100 μm was tested using 30 mm/s inlet velocity. By designing a significant smaller suspended microcantilever, using a geometry similar to the B1 version, better sensitivity can be achieved using lower velocities. This features set the proposed biosensor among the best candidate components for PoC or LoC devices for medical diagnostics and monitoring in terms of hematocrit measurement.

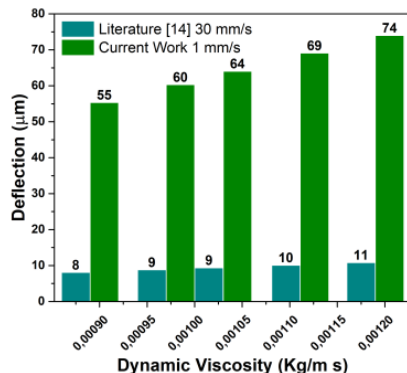


Figure 8. Comparison between the literature's data [14] and our work.

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